Chapter One

1. $AC\{D,B\} = ACDB + ACBD$, $A\{C,B\}D = ACBD + ABCD$, $C\{D,A\}B = CDAB + CADB$, and $\{C,A\}DB = CADB + ACDB$. Therefore $-AC\{D,B\} + A\{C,B\}D - C\{D,A\}B + \{C,A\}DB = -ACDB + ABCD - CDAB + ACDB = ABCD - CDAB = [AB,CD]$

In preparing this solution manual, I have realized that problems 2 and 3 in are misplaced in this chapter. They belong in Chapter Three. The Pauli matrices are not even defined in Chapter One, nor is the math used in previous solution manual. – Jim Napolitano

2. (a) $\text{Tr}(X) = a_0 \text{Tr}(1) + \sum_{\ell} \text{Tr}(\sigma_\ell) a_\ell = 2a_0$ since $\text{Tr}(\sigma_\ell) = 0$. Also $\text{Tr}(\sigma_k X) = a_0 \text{Tr}(\sigma_k) + \sum_{\ell} \text{Tr}(\sigma_k \sigma_\ell) a_\ell = \frac{1}{2} \sum_{\ell} \text{Tr}(\sigma_k \sigma_\ell + \sigma_\ell \sigma_k) a_\ell = \sum_{\ell} \delta_{k\ell} \text{Tr}(1) a_\ell = 2a_k$. So, $a_0 = \frac{1}{2} \text{Tr}(X)$ and $a_k = \frac{1}{2} \text{Tr}(\sigma_k X)$. (b) Just do the algebra to find $a_0 = (X_{11} + X_{22})/2$, $a_1 = (X_{12} + X_{21})/2$, $a_2 = i(-X_{21} + X_{12})/2$, and $a_3 = (X_{11} - X_{22})/2$.

3. Since $det(\boldsymbol{\sigma} \cdot \mathbf{a}) = -a_z^2 - (a_x^2 + a_y^2) = -|\mathbf{a}|^2$, the cognoscenti realize that this problem really has to do with rotation operators. From this result, and (3.2.44), we write

$$\det \left[\exp \left(\pm \frac{i \boldsymbol{\sigma} \cdot \hat{\mathbf{n}} \boldsymbol{\phi}}{2} \right) \right] = \cos \left(\frac{\boldsymbol{\phi}}{2} \right) \pm i \sin \left(\frac{\boldsymbol{\phi}}{2} \right)$$

and multiplying out determinants makes it clear that $det(\sigma \cdot \mathbf{a}') = det(\sigma \cdot \mathbf{a})$. Similarly, use (3.2.44) to explicitly write out the matrix $\sigma \cdot \mathbf{a}'$ and equate the elements to those of $\sigma \cdot \mathbf{a}$. With $\hat{\mathbf{n}}$ in the z-direction, it is clear that we have just performed a rotation (of the spin vector) through the angle ϕ .

4. (a) $\text{Tr}(XY) \equiv \sum_a \langle a|XY|a \rangle = \sum_b \sum_b \langle a|X|b \rangle \langle b|Y|a \rangle$ by inserting the identity operator. Then commute and reverse, so $\text{Tr}(XY) = \sum_b \sum_a \langle b|Y|a \rangle \langle a|X|b \rangle = \sum_b \langle b|YX|b \rangle = \text{Tr}(YX)$. (b) $XY|a \rangle = X[Y|a \rangle$] is dual to $\langle a|(XY)^\dagger$, but $Y|a \rangle \equiv |\beta \rangle$ is dual to $\langle a|Y^\dagger \equiv |\beta \rangle$ and $X[\beta]$ is dual to $\langle \beta|X^\dagger$ so that $X[Y|a \rangle]$ is dual to $\langle a|Y^\dagger X^\dagger$. Therefore $(XY)^\dagger = Y^\dagger X^\dagger$. (c) $\exp[if(A)] = \sum_a \exp[if(A)]|a \rangle \langle a| = \sum_a \exp[if(a)]|a \rangle \langle a|$ (d) $\sum_a \psi_a^*(\mathbf{x}')\psi_a(\mathbf{x}'') = \sum_a \langle \mathbf{x}'|a \rangle^* \langle \mathbf{x}''|a \rangle \langle a|\mathbf{x}'' \rangle = \langle \mathbf{x}''|\mathbf{x}' \rangle = \delta(\mathbf{x}'' - \mathbf{x}')$

5. For basis kets $|a_i\rangle$, matrix elements of $X \equiv |\alpha\rangle\langle\beta|$ are $X_{ij} = \langle a_i|\alpha\rangle\langle\beta|a_j\rangle = \langle a_i|\alpha\rangle\langle a_j|\beta\rangle^*$. For spin-1/2 in the $|\pm z\rangle$ basis, $\langle +|S_z=\hbar/2\rangle = 1$, $\langle -|S_z=\hbar/2\rangle = 0$, and, using (1.4.17a), $\langle \pm|S_x=\hbar/2\rangle = 1/\sqrt{2}$. Therefore

$$|S_z = \hbar/2\rangle\langle S_x = \hbar/2| \doteq \frac{1}{\sqrt{2}}\begin{pmatrix} 1 & 1 \\ 0 & 0 \end{pmatrix}$$

A[|i⟩ + |j⟩] = a_i|i⟩ + a_j|j⟩ ≠ [|i⟩ + |j⟩] so in general it is not an eigenvector, unless a_i = a_j.
That is, |i⟩ + |j⟩ is not an eigenvector of A unless the eigenvalues are degenerate.

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